

Physics-biophysics 1

Flow of fluids

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What is a fluid?

- *'something that can flow'*
- *'something that has no defined shape'*
- **fluid:** a substance that cannot support shear stress
- deformation:
 - *tensile stress (tension)*
 - *'tends to change the dimensions but not the shape'*
 - force \perp surface
 - $p = \frac{F}{A}$
 - *shear stress*
 - *'tends to change the shape but not the dimensions'*
 - force \parallel surface
 - $\tau = \frac{F}{A}$
- *fluid* is a common name for *liquids* and *gases* — in most cases, they can be treated the same way

Classification of fluids

1 Compressibility

- A **compressible:** its density (ρ) can change (eg, gases)
- B **incompressible:** its density (ρ) is constant (eg, liquids)

2 Internal friction (viscosity)

- A **viscous:** there IS internal friction
- B **non-viscous:** internal friction is negligible

- **ideal fluid:** incompressible and non-viscous

Basic quantities

- **current:** the mass or volume flowing through a cross-section of a tube in unit time

$$I := \frac{dm}{dt} \quad \text{or} \quad \frac{dV}{dt},$$

where m denotes the mass, t is the time and V stands for the volume

- the greater the cross-section area, the greater the current \rightarrow
- **current density:** the current flowing through a unit cross-section area

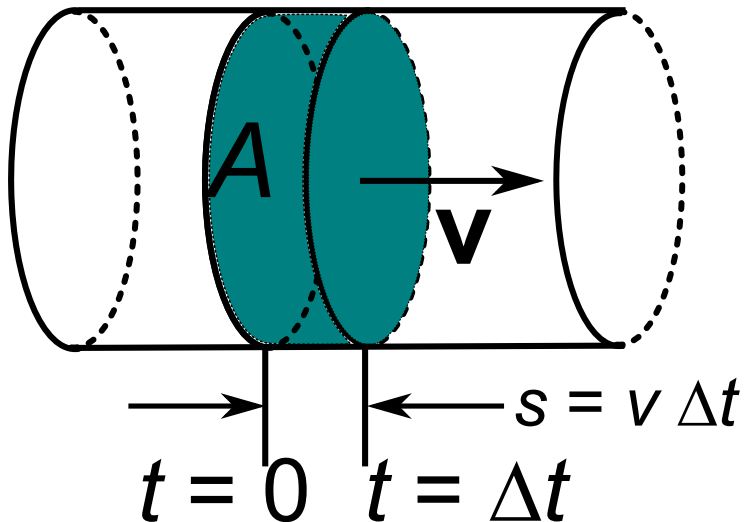
$$J := \frac{dI}{dA},$$

where A denotes the cross-section area

Flow of incompressible fluids

- liquids: generally incompressible
- gases: generally compressible
- even for gases, if the flow speed is not too high (< 50 m/s), ρ does not change significantly and the *flow* can be considered incompressible
- time dependence of the flow
 - **stationary (steady)**: speed and current are *independent of time*
 - **non-stationary (unsteady)**: speed and current *do depend on time*

Current and flow speed



Current and flow speed

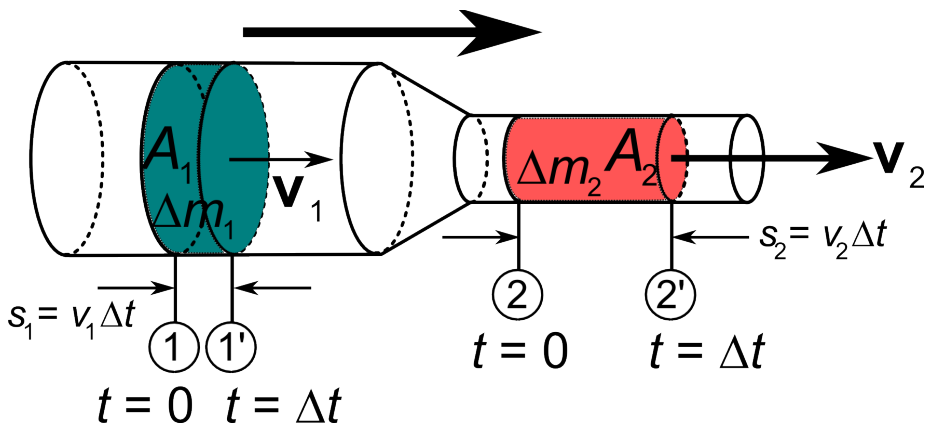
- let us assume that the fluid particles move parallel to each other with the same flow velocity \mathbf{v}
- this way, in time Δt they travel a distance $s = v\Delta t$
- from the perspective of fluid flow, this means that volume ΔV has been carried through a given cross-section area A of the tube, where

$$\Delta V = As = Av\Delta t$$

- thus the volume current of the flow is

$$I = \frac{\Delta V}{\Delta t} = \frac{Av\Delta t}{\Delta t} = Av$$

Illustration



Notations

- (1) and (2): selected discs in the fluid at time $t = 0$
- (1') and (2'): locations of the selected discs in at time $t = \Delta t$
- \mathbf{v}_1 : the velocity of the fluid at (1) and (1')
- \mathbf{v}_2 : the velocity of the fluid at (2) and (2')
- A_1 : area of the cross-section at (1) and (1')
- A_2 : area of the cross-section at (2) and (2')
- s_1 : distance travelled by the fluid between locations (1) and (1')
- s_2 : distance travelled by the fluid between locations (2) and (2')

Equation of continuity

- the fluid is incompressible, so
 - the density is the same between (1) and (1') and (2) and (2'): $\rho_1 = \rho_2 = \rho$
 - the mass flowing in at (1) is equal to the mass flowing out at (2):

$$\Delta m_1 = \Delta m_2 = \Delta m$$

$$\bullet \Delta m = \rho \Delta V_1 = \rho \Delta V_2 \rightarrow$$

- the volume flowing in at (1) is equal to the volume flowing out at (2):

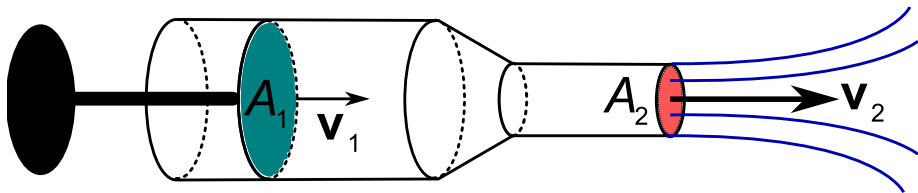
$$\Delta V_1 = \Delta V_2 = \Delta V$$

- $\Delta V_1 = A_1 s_1 = A_1 v_1 \Delta t$
- $\Delta V_2 = A_2 s_2 = A_2 v_2 \Delta t$
- $\Delta V_1 = \Delta V_2 \Rightarrow A_1 v_1 \Delta t = A_2 v_2 \Delta t$

$$A_1 v_1 = A_2 v_2$$

- since (1) and (2) were chosen arbitrarily, this must hold to any two cross-sections along the flow: $A \cdot v = \text{const}$
- this is the **equation of continuity**
- since $Av = I$, this means that the **volume current is constant along the tube**

Example: a syringe



$$A_1 v_1 = A_2 v_2$$

$$A_1 \gg A_2 \Rightarrow v_2 \gg v_1$$

Example: flow speed in blood vessels

Blood vessel	cross-section area [cm²]	velocity [cm/s]
Aorta	4.5	40
Arteries	20	9
Arterioles	400	0.45
Capillaries	4500	0.04
Veins	40	4.5
<i>Vena cava</i>	18	10

Homework: what is common to all the rows?

Flow of ideal fluids

- ideal fluid: incompressible and **non-viscous**
- no internal friction → no loss of mechanical energy
- mechanical energy: potential energy + kinetic energy

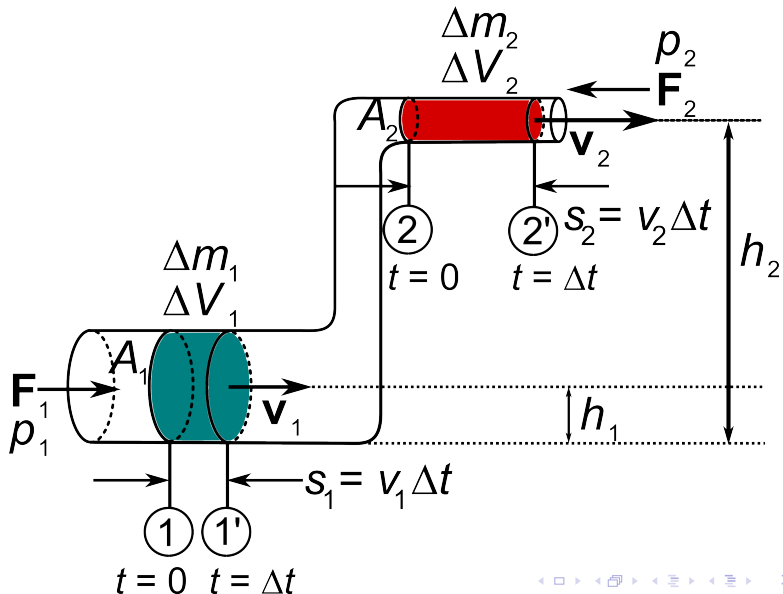
$$E = E_P + E_K$$

- the *conservation of mechanical energy* applies: external work done on the system = change in the mechanical energy of the system

$$W = \Delta E = \Delta E_P + \Delta E_K$$

- Bernoulli's law: a special form of the conservation of mechanical energy for the flow of ideal fluids

Illustration

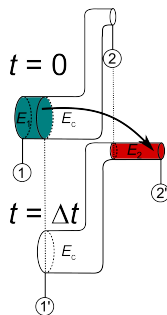


Notations

- (1) and (2): selected discs in the fluid at time $t = 0$
- (1') and (2'): locations of the selected discs in at time $t = \Delta t$
- \mathbf{v}_i : the velocity of the fluid at (i) and (i') ($i \in \{1, 2\}$)
- A_i : area of the cross-section at (i) and (i') ($i \in \{1, 2\}$)
- s_i : distance travelled by the fluid between locations (i) and (i') ($i \in \{1, 2\}$)
- h_i : height of the centre of mass of the fluid at (i) and (i') (as compared to an arbitrary reference level, $i \in \{1, 2\}$)
- p_i : pressure exerted by the rest of the fluid at (i) and (i') ($i \in \{1, 2\}$)
- \mathbf{F}_i : force exerted by the rest of the fluid at (i) and (i') ($i \in \{1, 2\}$)
- Δm_i : mass of the fluid between locations (i) and (i') ($i \in \{1, 2\}$)
- ΔV_i : volume of the fluid between locations (i) and (i') ($i \in \{1, 2\}$)
- ρ_i : density of the fluid between locations (i) and (i') ($i \in \{1, 2\}$)

Energy changes

- how does the energy of the fluid between (1) and (2) change?
- kinetic and potential energies only depend on the position
- E_C : the energy of the part of the fluid which does not change between $t = 0$ and $t = \Delta t$
- $E = E_1 + E_C$
- $E' = E_C + E_2$
- $\Delta E = E' - E = E_2 - E_1$
- only the energy of the marked sections changes



Implications of the equation of continuity

$$\rho_1 = \rho_2 = \rho$$

$$\Delta m_1 = \Delta m_2 = \Delta m$$

$$\Delta V_1 = \frac{\Delta m_1}{\rho} = \frac{\Delta m_2}{\rho} = \Delta V_2$$

$$\Delta V_1 = A_1 s_1 = A_1 v_1 \Delta t$$

$$\Delta V_2 = A_2 s_2 = A_2 v_2 \Delta t$$

Difference between the fluid states at $t = 0$ and $t = \Delta t$: mass Δm is transported from (1) to (2)

What happens to this mass?

- 1 the rest of the fluid does work on it
- 2 moved from h_1 to $h_2 \rightarrow$ its potential energy changes
- 3 its velocity changes \rightarrow its kinetic energy changes

Work done by the rest of the fluid

- definition of pressure:

$$p = \frac{F}{A}$$

$$F = pA$$

- work done by the fluid at (1) and (2):

$$W_1 = F_1 s_1 = F_1 v_1 \Delta t = p_1 A_1 v_1 \Delta t = p_1 \Delta V = p_1 \frac{\Delta m}{\rho}$$

$$W_2 = -F_2 s_2 = -F_2 v_2 \Delta t = -p_2 A_2 v_2 \Delta t = -p_2 \Delta V = -p_2 \frac{\Delta m}{\rho}$$

- negative sign in W_2 : the direction of the force is opposite to that of the displacement
- total external work:

$$W = W_1 + W_2 = \frac{\Delta m}{\rho} (p_1 - p_2)$$

Potential and kinetic energies

Gravitational potential energy: $E_P = mgh$

- at (1): $E_{P,1} = \Delta mgh_1$
- at (2): $E_{P,2} = \Delta mgh_2$
- change: $\Delta E_P = E_{P,2} - E_{P,1} = \Delta mg(h_2 - h_1)$

Kinetic energy: $E_K = \frac{1}{2}mv^2$

- at (1): $E_{K,1} = \frac{1}{2}\Delta mv_1^2$
- at (2): $E_{K,2} = \frac{1}{2}\Delta mv_2^2$
- change: $\Delta E_K = E_{K,2} - E_{K,1} = \frac{1}{2}\Delta m(v_2^2 - v_1^2)$

Bernoulli's equation

- conservation of mechanical energy: $W = \Delta E = \Delta E_P + \Delta E_K$

$$\frac{\Delta m}{\rho} (p_1 - p_2) = \Delta mg(h_2 - h_1) + \frac{1}{2} \Delta m (v_2^2 - v_1^2) \quad // \cdot \frac{\rho}{\Delta m}$$

$$p_1 - p_2 = \rho gh_2 - \rho gh_1 + \frac{1}{2} \rho v_2^2 - \frac{1}{2} \rho v_1^2$$

$$p_1 + \rho gh_1 + \frac{1}{2} \rho v_1^2 = p_2 + \rho gh_2 + \frac{1}{2} \rho v_2^2$$

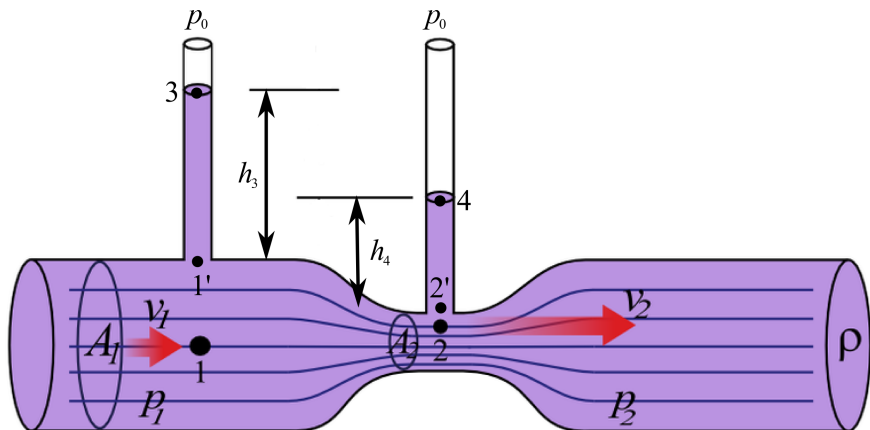
- fluid discs (1) and (2) were chosen arbitrarily, so this must hold to any two cross-sections along the flow: $p + \rho gh + \frac{1}{2} \rho v^2 = \text{const}$

- this is **Bernoulli's law**

Interpreting Bernoulli's law

- p : **static pressure**
- ρgh : **hydrostatic pressure**
- $\frac{1}{2}\rho v^2$: **dynamic pressure**
- **total pressure** = static pressure + hydrostatic pressure + dynamic pressure
- Bernoulli's law in other words: *the total pressure is constant along the tube*

Example: Venturi tube



Example: Venturi tube

- from the equation of continuity:

$$A_1 v_1 = A_2 v_2$$

$$A_1 \gg A_2 \Rightarrow v_2 \gg v_1$$

- apply Bernoulli's law to compare 1 and 2:

$$p_1 + 0 + \frac{1}{2}\rho v_1^2 = p_2 + 0 + \frac{1}{2}\rho v_2^2$$

- since $v_2 \gg v_1 \Rightarrow p_1 \gg p_2$
- 1 and 3 cannot be compared, because they are in different tubes; but 1' and 3 can (1' is at the beginning of the vertical tube)
- the static pressures at 1 and 1' are the same: $p_{1'} = p_1$
- the dynamic pressure at 1' and 3 is 0, because the fluid does not flow in the vertical tubes
- apply Bernoulli's law to compare 1' and 3 (the reference level is now at 1'):

$$p_1 + 0 + 0 = p_0 + \rho h_3 g + 0,$$

where p_0 is the atmospheric pressure

Example: Venturi tube

- 2 and 4 cannot be compared, because they are in different tubes; but 2' and 4 can (2' is at the beginning of the vertical tube)
- the static pressures at 2 and 2' are the same: $p_{2'} = p_2$
- the dynamic pressure at 2' and 4 is 0, because the fluid does not flow in the vertical tubes
- apply Bernoulli's law to compare 2' and 4 (the reference level is now at 2')

$$p_2 + 0 + 0 = p_0 + \rho h_4 g + 0$$

- if we compare h_3 and h_4

$$h_3 = \frac{p_1 - p_0}{\rho g} \gg h_4 = \frac{p_2 - p_0}{\rho g},$$

because we have seen that $p_1 \gg p_2$

- the Venturi tube proves Bernoulli's law — *at the wider section of the tube, where flow speed is smaller, the static pressure is higher as compared to narrower sections of the tube*
- the role of vertical tubes: to make the differences in static pressure visible

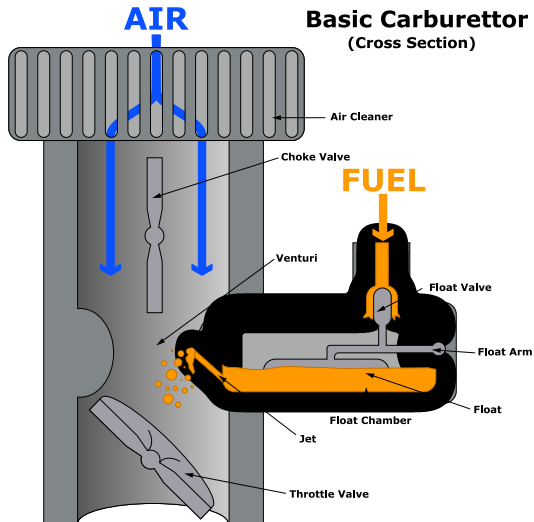
Example: Bunsen burner

- flow speed in the narrower section is greater \Rightarrow lower static pressure
- static pressure within the tube is less than the atmospheric pressure
- as a result, air flows into the tube
- the air influx feeds the flame at the top



Example: carburettor

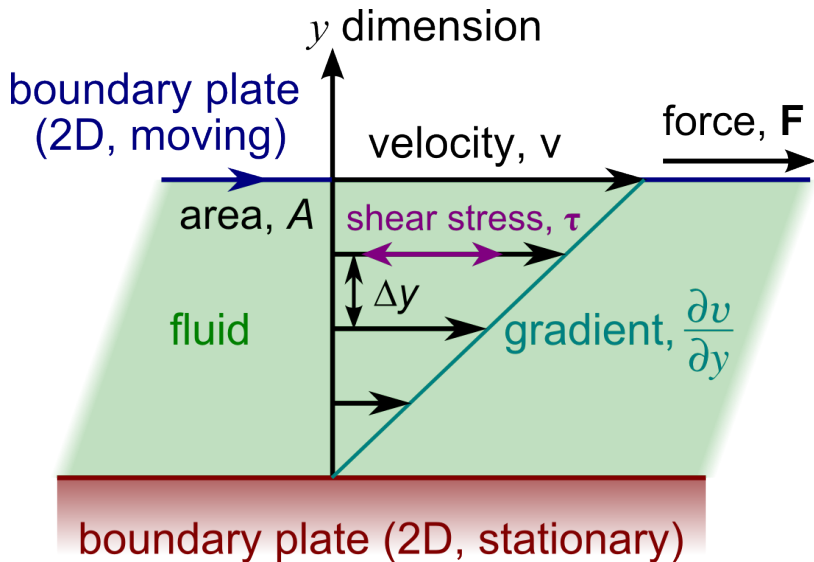
- air injected into the carburettor chamber
- in the narrow throat, the air is moving at its fastest speed and therefore it is at its lowest pressure
- low pressure in the chamber pumps the fuel into the chamber, where it is mixed with air



What is viscosity?

- fluids (by definition) cannot *support* shearing stress
- but they *do resist* shearing motion
- interactions between the particles of the fluid \Rightarrow friction
- force is needed to make fluid layers slide upon each other
- viscosity: the resistance of a fluid to shearing motion due to internal friction
- everyday term for viscosity: 'thickness' (riddle: which English saying involves viscosity?)
- etymology: < Latin *viscum* 'mistletoe' (mistletoe glue was used to catch birds)

Illustration



What happens in a viscous fluid?

- two parallel solid plates with surface area A with viscous fluid between them
- lower plate: fixed
- upper plate: pulled sideways with force \mathbf{F}
- bottom layers of fluid: stick to the lower plate due to adhesion, velocity: $v = 0$
- top layers of fluid: stick to the upper plate due to adhesion, velocity: $v \neq 0$
- we imagine fluid flow as fluid layers sliding upon each other
- particles within the same layer move together with the same velocity: $\mathbf{v}(y)$

What happens in a viscous fluid?

- thickness of layers: Δy
- speed difference between neighbouring layers: Δv
- instead of the force, it is more practical to use the shear stress τ

$$\tau := \frac{F}{A}$$

- how much the velocity changes across layers — *velocity gradient*:

$$\frac{\Delta v}{\Delta y}$$

- in reality, there are no homogeneous layers, the velocity changes continuously across the cross-section of the tube
- \Rightarrow layer thickness Δy should be made infinitely small
- thus the velocity gradient (also called the **shear rate**):

$$\lim_{\Delta y \rightarrow 0} \frac{\Delta v}{\Delta y} = \frac{dv}{dy}$$

Newton's law of friction

- the stronger the internal friction, the more the layers move together \Rightarrow the less the velocity difference between layers (shear rate)
- non-viscous fluid (eg, water): we can slide a solid plate on top of it without making the whole bulk of fluid moving
- viscous fluid (eg, honey): sliding a solid plate on its surface makes the whole bulk of fluid move
- in many fluids, shear rate $\frac{dv}{dy}$ is proportional to the shear stress τ
- the constant of proportionality is called the **viscosity** η
- **Newton's law of friction:**

$$\tau = \frac{F}{A} = \eta \frac{dv}{dy}$$

- **Newtonian fluids:** fluids which obey Newton's law of friction

Viscosity

- using Newton's law of friction, we can define viscosity:

$$\eta := \tau \left(\frac{dv}{dy} \right)^{-1}$$

- the unit of viscosity:

$$[\eta] = 1 \text{ Pa} \cdot \text{s} \quad \left(\leftarrow 1 \frac{\text{N}}{\text{m}^2} \div \frac{\text{m/s}}{\text{m}} = 1 \frac{\text{Pa}}{1/\text{s}} \right)$$

Fluid	glycerine	blood	water	air
Viscosity [Pa · s]	0.83	0.02–0.04	0.001	0.00001

Table 1: Typical viscosity values

Viscosity as a diagnostic tool

- blood viscosity measurement has a great potential as a diagnostic tool
- viscosity changes in conditions of
 - myocardial infarction
 - coronary occlusion
 - arteriosclerosis
 - diabetes, &c
- it is not yet known whether viscosity changes are just symptoms or themselves contribute to the disorders they accompany

Molecular origin of viscosity in gases

- in ideal gases, molecules only interact through collisions
- source of friction: *momentum exchange between layers*
- molecules in gas flow:
 - ordered motion — due to pressure difference
 - disordered motion — due to temperature (usually with much greater speed than ordered motion)
- disordered motion \Rightarrow molecules may enter other layers, and through collisions, change the momentum of the layer
- momentum: $\mathbf{p} := m\mathbf{v}$; Newton's 2nd law: $\mathbf{F} = m\mathbf{a} = \frac{d\mathbf{p}}{dt}$
- momentum exchange \Rightarrow force of friction between layers
- force of friction \Rightarrow viscosity

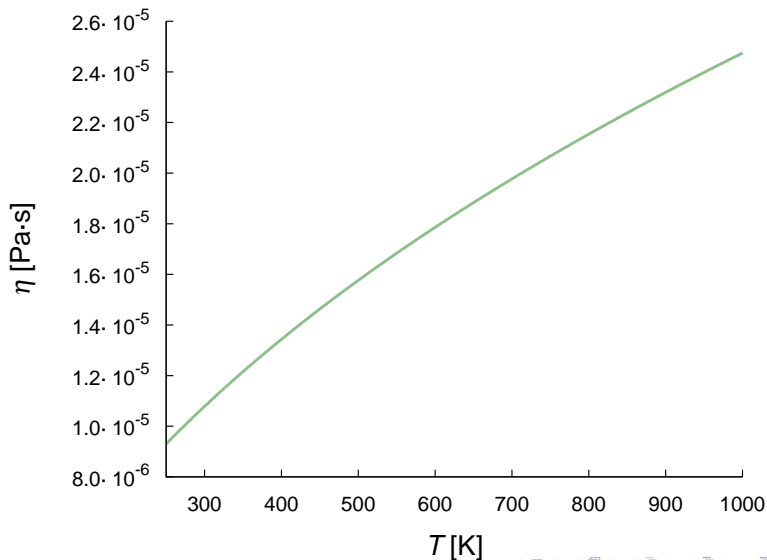
Temperature dependence of viscosity in gases

- equipartition theory for ideal gases: $\frac{1}{2}mv^2 = \frac{f}{2}kT$, where v is the average speed of gas molecules, k is Boltzmann's constant, T is the thermodynamic temperature (in Kelvins) and f denotes the degrees of freedom
- higher temperature \Rightarrow faster disordered motion
- faster disordered motion \Rightarrow more collisions in unit time \Rightarrow stronger force of friction
- **the viscosity of gases increases with temperature**

$$\eta = \eta_0 \sqrt{\frac{T}{T_0}} \cdot \frac{1 + \frac{C}{T_0}}{1 + \frac{C}{T}}$$

where η_0 is the viscosity of the gas at absolute temperature T_0 and C is called the Sutherland constant

Temperature dependence of viscosity in gases



Molecular origin of viscosity in liquids

- in liquids, there are strong interactions between molecules
- attractive forces \Rightarrow internal friction
- flow requires free space in order for the molecules to move and change places
- *How can liquids flow at all when there is not enough free space within them?*

Frenkel's theory

- thermal motion results in occasional molecular gaps in the liquid
- it provides room for liquid molecules to move \Rightarrow fluid flow
- the higher the concentration of the gaps, the more easily liquid molecules can move \Rightarrow lower viscosity

Temperature dependence of viscosity in liquids

- liquid molecules possess potential energy due to interactions with neighbouring molecules
- to create a gap, a molecule must be freed \Rightarrow some activation energy E must be invested
- source of activation energy: kinetic energy of liquid molecules
- *How many molecules have kinetic energy greater than or equal to E ?*
- Boltzmann distribution \Rightarrow

$$N \propto T \cdot e^{-\frac{E}{kT}},$$

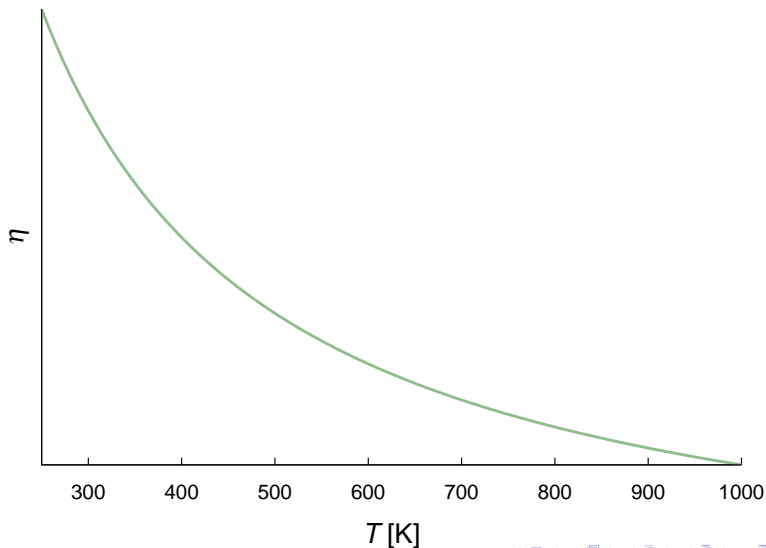
where N is the number of molecules above the activation energy E , k is the Boltzmann constant and T is the absolute temperature

- the higher the concentration of the gaps, the lower the viscosity $\Rightarrow \eta \propto \frac{1}{N}$:

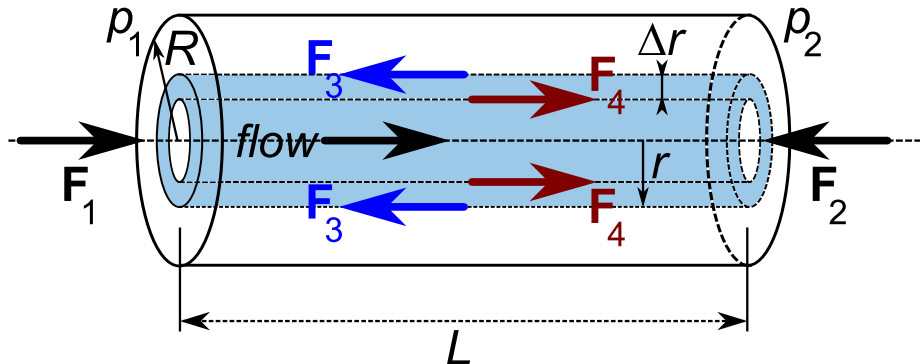
$$\eta \propto \frac{1}{T} e^{\frac{E}{kT}}$$

- **the viscosity of liquids decreases with temperature**

Temperature dependence of viscosity in liquids



Hagen–Poiseuille law: illustration



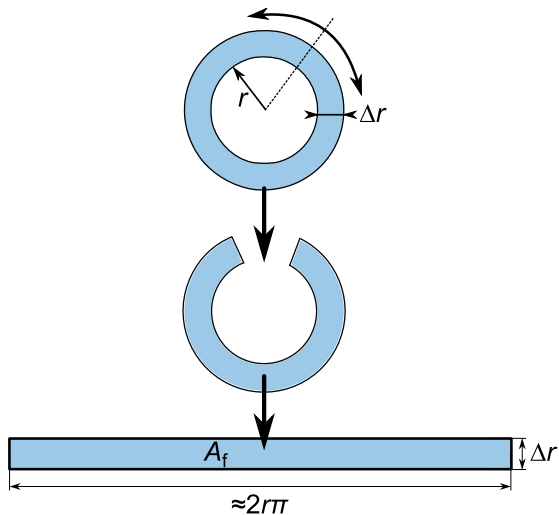
Notations

- p_1 : pressure exerted by the rest of the fluid at the beginning of the tube
- p_2 : pressure exerted by the rest of the fluid at the end of the tube
- F_1 : propelling force due to fluid pressure at the beginning of the tube
- F_2 : blocking force due to fluid pressure at the end of the tube
- F_3 : force of friction slowing down the selected fluid volume along the outer surface
- F_4 : force of friction accelerating the selected fluid volume along the inner surface
- Δr : thickness of the selected fluid volume
- r : distance from the axis
- R : radius of the tube
- A_f : area of the front surface of the selected volume
- $A(r)$: lateral surface area (area of the side) of the selected volume (different for inner and outer surfaces)

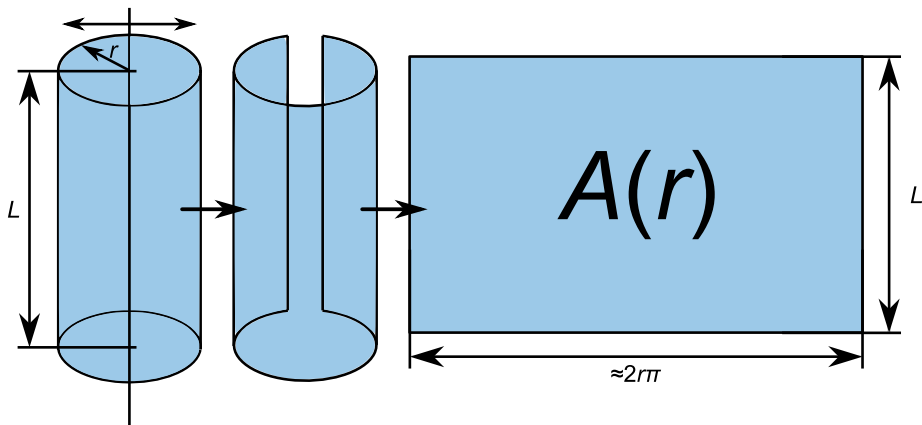
Properties of the flow

- the tube has a cylindrical symmetry \Rightarrow fluid layers will have a cylindrical symmetry
- particles at the same distance r from the axis travel with the same velocity
- particles at the wall: do not move due to adhesion to the wall ($v = 0$)
- particles in the axis: maximum speed
- we select a hollow tube of fluid with thickness Δr at distance r from the axis
- fluid moving *within* the selected volume: moving faster \Rightarrow accelerating the selected fluid volume due to friction
- fluid moving *outside* the selected volume: moving slower \Rightarrow decelerating the selected fluid volume due to friction
- pressure of the rest of the fluid: accelerates the selected volume at one end, decelerates it at the other
- we consider a stationary flow \Rightarrow acceleration is zero \Rightarrow the vector sum of all the forces acting on the selected fluid volume is zero

Front surface area



Lateral surface area



Forces acting on the selected fluid volume

Forces due to pressure

- pressure at the beginning is greater than pressure at the end: $p_1 > p_2 \Rightarrow$ direction of the flow
- $A_f \approx 2r\pi\Delta r$ (see the figure)
- $F_1 = p_1 A_f = p_1 2r\pi\Delta r$
- $F_2 = -p_2 A_f = -p_2 2r\pi\Delta r$

Forces due to viscosity

- Newton's law of friction: $\tau = \eta \frac{dv}{dr} \Rightarrow F = A\eta \frac{dv}{dr}$
- $\frac{dv}{dr}$ is negative because the velocity decreases with r
- lateral surface area at distance r from the axis: $A(r) = L \cdot 2r\pi$ (see the figure)
- force of outer layer (blocking): $F_3 = \eta A(r + \Delta r) \left. \frac{dv}{dr} \right|_{r+\Delta r}$
- force of inner layer (propelling): $F_4 = -\eta A(r) \left. \frac{dv}{dr} \right|_r$

Obtaining the velocity

- we consider a stationary flow \Rightarrow acceleration is zero \Rightarrow the vector sum of all the forces acting on the selected fluid volume is zero (Newton's second law)

$$\sum F = F_1 + F_2 + F_3 + F_4 = 0$$

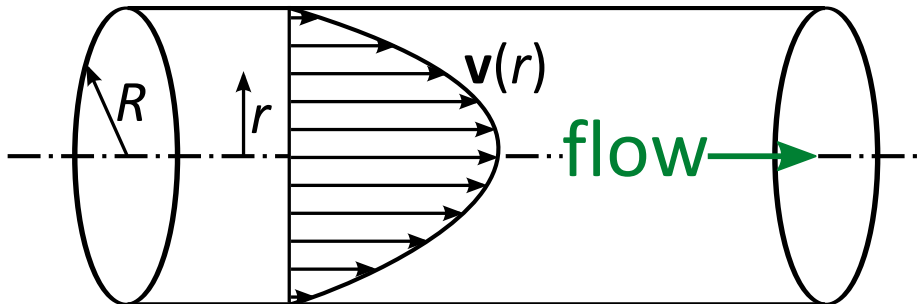
$$p_1 2r\pi\Delta r - p_2 2r\pi\Delta r + \eta A(r + \Delta r) \left. \frac{dv}{dr} \right|_{r+\Delta r} - \eta A(r) \left. \frac{dv}{dr} \right|_r = 0$$

- solving this differential equation yields

$$v(r) = \frac{p_1 - p_2}{4\eta L} (R^2 - r^2)$$

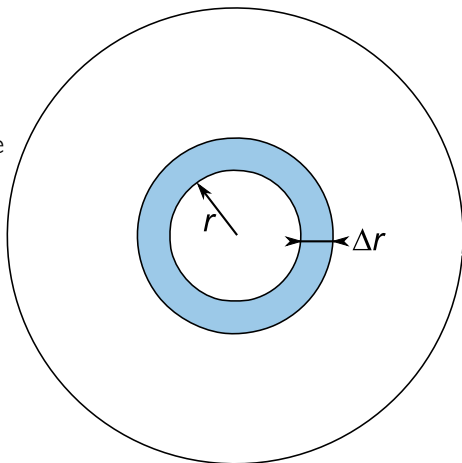
- it is called the **parabolic velocity profile**
- the farther we are from the axis, the less the velocity
- $v(0) = \frac{p_1 - p_2}{4\eta L} R^2$ — maximum value
- $v(R) = 0$

Parabolic velocity profile

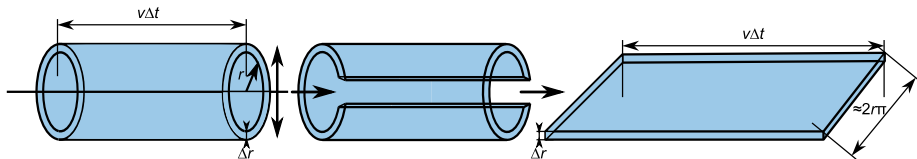


Obtaining the current

- we now know the velocity, but not the current
- let us consider the *current* flowing through a Δr thick ring at distance r from the axis: $I(r)$
- the ring is very narrow $\Rightarrow v(r)$ can be considered as constant within the ring
- here we consider the *volume current*: $I(r) = \frac{\Delta V}{\Delta t}$, where ΔV is the fluid volume flowing through our selected ring in time Δt



Obtaining the volume ΔV



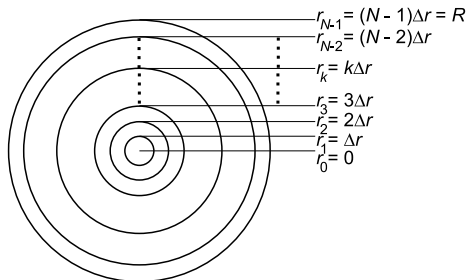
$$\Delta V = 2r\pi \cdot v(r)\Delta t \cdot \Delta r$$

$$l(r) = \frac{\Delta V}{\Delta t} = \frac{2r\pi \cdot v(r)\Delta t \cdot \Delta r}{\Delta t} = 2r\pi \cdot v(r) \cdot \Delta r$$

Summing the currents

- to obtain the total current I , we have to add up the currents $I(r)$ of all the rings constituting the total cross section
- the cross section is broken down into N rings of equal thickness Δr
- for the k^{th} ring, the inner radius is $r_k := k\Delta r$
- the current through the k^{th} ring:
 $I_k := I(r_k) = 2r_k\pi \cdot v(r_k) \cdot \Delta r$
- the total current:

$$I \approx \sum_{k=0}^{N-1} I_k = \sum_{k=0}^{N-1} 2r_k\pi \cdot v(r_k) \cdot \Delta r$$



Hagen–Poiseuille equation

- exact solution obtained through integration

$$I = \lim_{\Delta r \rightarrow 0} \sum_{k=0}^{N-1} 2r_k \pi \cdot v(r_k) \cdot \Delta r = \int_0^R 2r \pi \cdot v(r) dr$$

$$I = \int_0^R 2r \pi \cdot v(r) dr = \int_0^R 2r \pi \cdot \frac{\Delta p}{4\eta L} (R^2 - r^2) dr = \frac{\pi}{2\eta L} \cdot \Delta p \cdot \int_0^R (rR^2 - r^3) dr$$

$$I = \frac{\pi}{2\eta L} \cdot \Delta p \cdot \left(R^2 \int_0^R r dr - \int_0^R r^3 dr \right) = \frac{\pi}{2\eta L} \cdot \Delta p \cdot \left(R^2 \left[\frac{r^2}{2} \right]_0^R - \left[\frac{r^4}{4} \right]_0^R \right)$$

$$I = \frac{\pi}{2\eta L} \cdot \Delta p \cdot \left\{ R^2 \left(\frac{R^2}{2} - 0 \right) - \left(\frac{R^4}{4} - 0 \right) \right\} = \frac{\pi}{2\eta L} \cdot \Delta p \frac{R^4}{4}$$

$$I = \frac{\pi}{8\eta L} \cdot \Delta p \cdot R^4$$

- this is the **Hagen–Poiseuille equation**

Medical examples

Muscles in use (eg, sports)

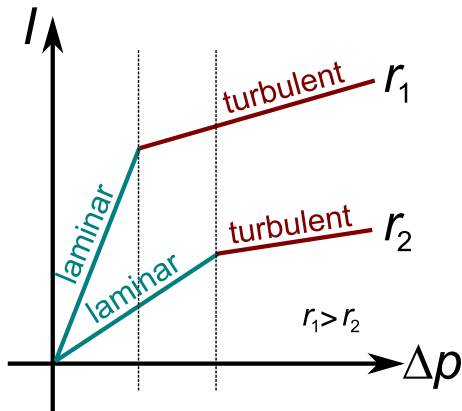
- oxygen demand increases
- more blood (higher current) needed to supply this oxygen
- blood vessels dilate (become wider)
- $I \propto R^4$
- the current through the dilated blood vessels increases drastically

Arteriosclerosis

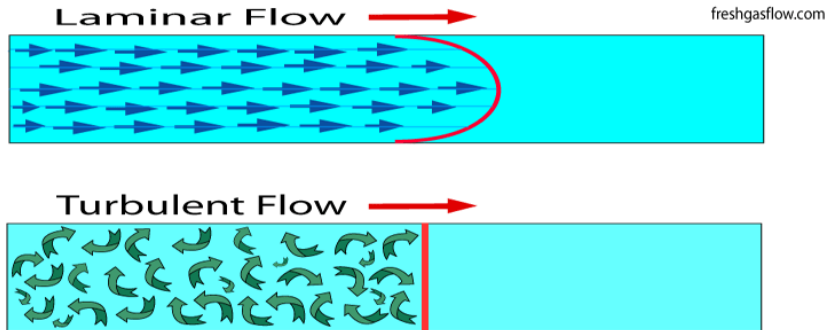
- blood vessels become rigid and *narrow*
- narrower tubes $\Rightarrow R$ drops
- $I \propto R^4 \Rightarrow I$ drops drastically
- much higher pressure needed to maintain the same current
- the heart may not sustain the increased load

Laminar and turbulent flow

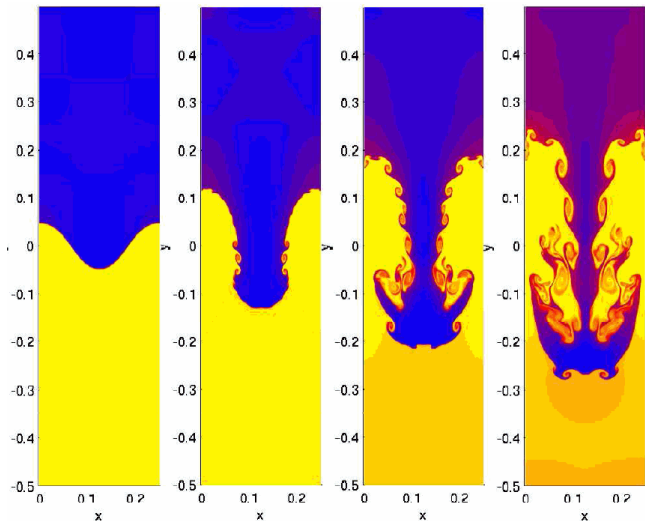
- **laminar flow:** a fluid flows in parallel layers, with no disruption between the layers
- **turbulent flow:** eddies form, flow properties become chaotic
- turbulent flow carries less volume for the same pressure difference: eddies reverse the direction of the flow and energy is dissipated



Laminar and turbulent flow



Laminar and turbulent flow



Laminar and turbulent flow



Laminar and turbulent flow



Reynolds number

- Is the flow laminar or turbulent? \Leftarrow *Reynolds number*

$$\text{Re} = \frac{\rho v L}{\eta}$$

- ρ : density of the fluid
- v : velocity of the flow
- L : characteristic length (eg, radius of the tube)
- η : viscosity of the fluid

- Re: dimensionless empirical number
- How is it used to predict turbulence? \Rightarrow *transition Reynolds numbers*
 - $\text{Re}_{\text{flow}} < \text{Re}_1$: laminar flow
 - $\text{Re}_1 \leq \text{Re}_{\text{flow}} < \text{Re}_2$: transition region — both laminar and turbulent flow are possible
 - $\text{Re}_2 \leq \text{Re}_{\text{flow}}$: turbulent flow
- transition Reynolds numbers are determined experimentally (eg, wind tunnels)
- transition Reynolds number values are similar for different fluids, they mostly depend on flow geometry

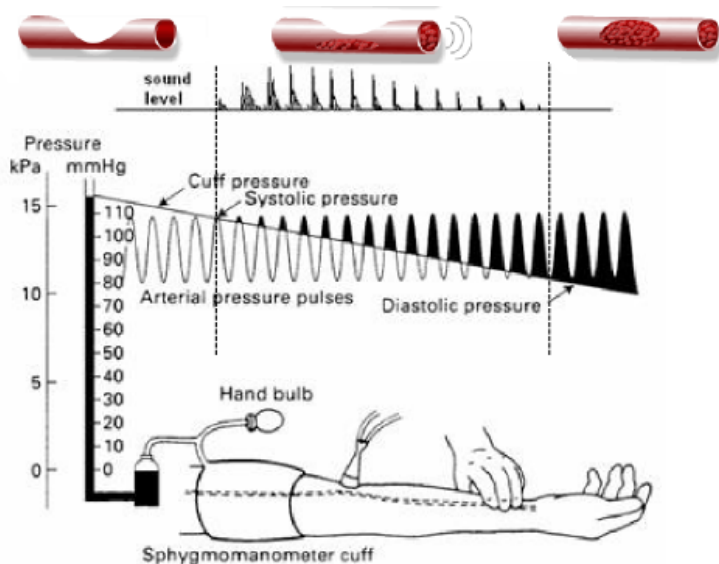
Example: blood pressure measurement



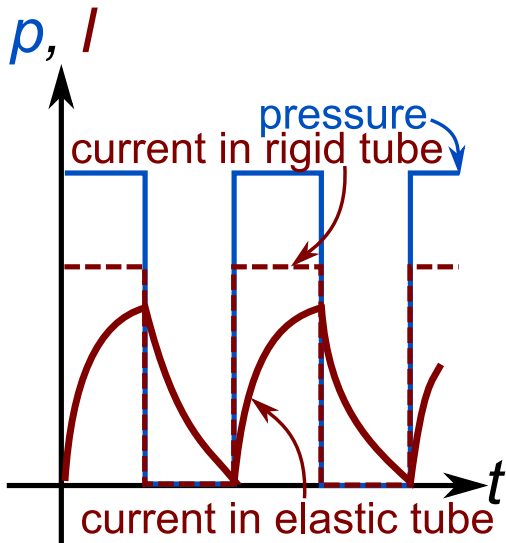
Example: blood pressure measurement

- the cuff (called the *sphygmomanometer cuff*) is inflated
- the cuff squeezes the artery \Rightarrow radius of the artery is reduced
- equation of continuity: flow velocity will increase through a narrower cross section
- increased velocity \Rightarrow Reynolds number increases \Rightarrow turbulent flow
- turbulent flow is noisy (*Korotkoff sounds*)
- these sounds are audible through the stethoscope

Example: blood pressure measurement



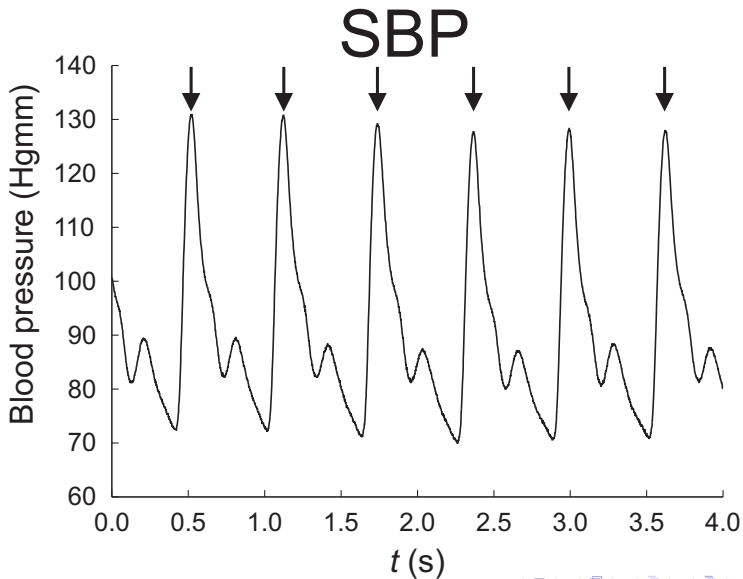
Pulsatile flow



Pulsatile flow

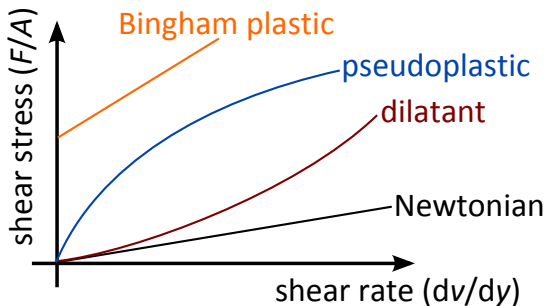
- eg, heart: operates in cycles of contraction and relaxation \Rightarrow *pulsatile flow*
- simple model: a tap opened and closed periodically
- elastic tube
 - high pressure: the walls of the tube expand and store elastic energy; the kinetic energy of the flow is converted to elastic energy
 - low pressure: the walls of the tube relax from their expanded state; their elastic energy is converted back to kinetic energy
 - this way, flow is maintained even when the pressure difference is low
 - physiological significance: the elasticity of veins makes blood pressure more stable
- rigid tube
 - no expansion or relaxation possible
 - lost mechanical energy is converted to heat

Blood pressure



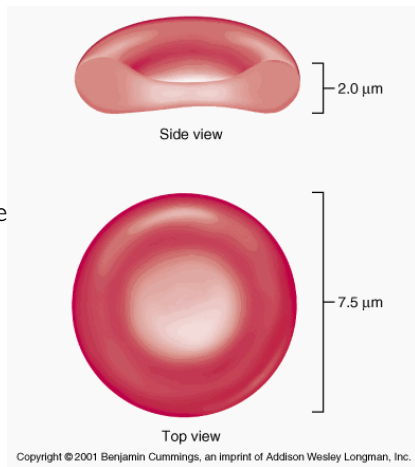
Types of non-Newtonian fluids

- Newtonian fluids
 - obey Newton's law of friction
 - their viscosity does not depend on shear rate
- non-Newtonian fluids
 - do *not* obey Newton's law of friction
 - their viscosity depends on shear rate, and in many cases, on the history of the fluid



Flow properties of blood

- volume of blood: 5–6 litres
- cellular components of blood:
 - erythrocytes (red blood cells)
 - leukocytes (white blood cells)
 - platelets
- the diameter of erythrocytes is greater than that of capillaries \Rightarrow erythrocytes can only squeeze through capillaries due to their elasticity
- **bolus flow**: erythrocytes block normal flow, blood plasma between them is forced to rotate \Rightarrow better circulation
- anomalies: sickle anaemia — reduced membrane elasticity of erythrocytes \Rightarrow being trapped (ageing of erythrocytes also reduces elasticity)



Bolus flow

